

## A POSITIVE SOLUTION OF A SCHRÖDINGER-POISSON SYSTEM WITH CRITICAL EXPONENT

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(Communicated by Irena Lasiecka)

### Abstract

We use variational methods to study the existence of at least one positive solution of the following Schrödinger-Poisson system

$$\begin{cases} -\Delta u + u + l(x)\phi u = k(x)|u|^{2^*-2}u + \mu h(x)|u|^{q-2}u & \text{in } \mathbb{R}^3, \\ -\Delta \phi = l(x)u^2 & \text{in } \mathbb{R}^3, \end{cases}$$

under some suitable conditions on the non-negative functions  $l, k, h$  and constant  $\mu > 0$ , where  $2 \leq q < 2^*$  (critical Sobolev exponent).

**AMS Subject Classification:** 35J20, 35J70

**Keywords:** Schrödinger-Poisson system; Variational methods; Critical growth; Positive solution

## 1 Introduction

In this paper, we study the existence of solutions of the system (1.2) involving a critical growth with the following form

$$\begin{cases} -\Delta u + u + l(x)\phi u = k(x)|u|^{2^*-2}u + \mu h(x)|u|^{q-2}u & \text{in } \mathbb{R}^3, \\ -\Delta \phi = l(x)u^2 & \text{in } \mathbb{R}^3, \end{cases} \quad (1.1)$$

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where  $2 \leq q < 2^*$ . We use the standard Mountain Pass Theorem to show the existence of a solution. However, since the nonlinearity involves a critical exponent, the Sobolev embedding  $H^1(\mathbb{R}^3) \hookrightarrow L^s(\mathbb{R}^3)$  ( $2 \leq s \leq 6$ ) is not compact. This will create great difficulties in the proof of the Palais-Smale condition. We will transform the problem into a nonlocal elliptic equation in  $\mathbb{R}^3$  and we also consider the limiting case  $q = 2$ .

It is known that the Schrödinger-Poisson systems have a strong physical meaning because they appear in quantum mechanics models (see e.g. [6, 9, 22]) and in semiconductor theory (see e.g. [4, 5, 23, 24]). In particular, systems like (1.2) have been introduced in Benci-Fortunato [4, 5] as a model describing solitary waves for the nonlinear stationary Schrödinger equations in three-dimensional space interacting with the electrostatic field which is not a priori assigned. Further applications to superconductors are currently under investigation.

Very recently, Cerami-Vaira [10] studied the existence of positive solutions for the Schrödinger-Poisson system

$$\begin{cases} -\Delta u + u + l(x)\phi u = f(x, u) & \text{in } \mathbb{R}^3, \\ -\Delta \phi = l(x)u^2 & \text{in } \mathbb{R}^3, \end{cases} \quad (1.2)$$

where they considered  $f(x, u) = k(x)|u|^{p-2}u$  with  $4 < p < 6$  and assumed that  $l \in L^2(\mathbb{R}^3)$  and  $k : \mathbb{R}^3 \rightarrow \mathbb{R}$  are non-negative functions satisfying  $\lim_{|x| \rightarrow +\infty} l(x) = 0$ ,  $l \not\equiv 0$ ,  $\lim_{|x| \rightarrow +\infty} k(x) = k_\infty > 0$  and  $k(x) - k_\infty \in L^{6/(6-p)}(\mathbb{R}^3)$ .

After Cerami-Vaira [10] many researchers have looked to problem (1.2), such as D'Avenia-Pomponio-Vaira [18], Li-Peng-Wang [21], Sun-Chen-Nieto [27] and Vaira [30], under various assumptions on the non-constant function  $l$ . Similar problems continue to attract attention as one can see from the latest works of He-Zou [20] and their references.

Before Cerami-Vaira [10] similar problems to (1.2), with constant function  $l$ , had also been widely investigated. We point out the works of Ambrosetti-Ruiz [2], Coclite [12], D'Avenia [17], D'Aprile et al. [13, 14, 15, 16], Ruiz [26] and others. Among of these, Azzollini-Pomponio [3], D'Aprile-Mugnai [14] and Zhao-Zhao [32] dealt with critical exponent case.

There are no existence results about system (1.1) with non-constant function  $l$ . In Zhao-Zhao [32], they studied a similar system to (1.1) with function  $l = 1$ . They established the existence of at least one positive solution for  $4 \leq q < 2^*$  and at least one positive radial solution for  $2 < q < 4$  with some restrictions on functions  $k$ ,  $h$  and  $\mu$ . Moreover, note that there was no information about the case where  $q = 2$ .

The main result, in this work, generalizes some of above results. We consider the following hypotheses ( $H$ ):

- ( $H_l$ )  $l \in L^2(\mathbb{R}^3) \cap L^\infty(\mathbb{R}^3)$ ,  $l(x) \geq 0$  for any  $x \in \mathbb{R}^3$  and  $l \not\equiv 0$ ;
- ( $H_{k_1}$ )  $k(x) \geq 0$  for any  $x \in \mathbb{R}^3$ ;
- ( $H_{k_2}$ ) There exists  $x_0 \in \mathbb{R}^3$ ,  $\delta_1 > 0$  and  $\rho_1 > 0$  such that  $k(x_0) = \max_{\mathbb{R}^3} k(x)$  and  $|k(x) - k(x_0)| \leq \delta_1|x - x_0|^\alpha$  for  $|x - x_0| < \rho_1$  with  $1 \leq \alpha < 3$ ;
- ( $H_{h_1}$ )  $h \in L^{6/(6-q)}(\mathbb{R}^3)$  and  $h(x) \geq 0$  for any  $x \in \mathbb{R}^3$  and  $h \not\equiv 0$ ;
- ( $H_{h_2}$ ) There are  $\delta_2 > 0$  and  $\rho_2 > 0$  such that  $h(x) \geq \delta_2|x - x_0|^{-\beta}$  for  $|x - x_0| < \rho_2$  and  $2 - \frac{q}{2} < \beta < 3$ , where  $x_0$  is given by ( $H_{k_2}$ );

$(H_{h_\mu})$   $0 < \mu < \bar{\mu}$  when  $2 \leq q < 4$ ;  $\mu > 0$  when  $4 \leq q < 6$ , where  $\bar{\mu}$  is defined by

$$\bar{\mu} := \mu_h = \inf_{u \in H^1(\mathbb{R}^3) \setminus \{0\}} \left\{ \int_{\mathbb{R}^3} (|\nabla u|^2 + u^2) dx : \int_{\mathbb{R}^3} h(x)|u|^q dx = 1 \right\}.$$

*Remark 1.1.* The hypotheses  $(H_{k_1})$  and  $(H_{k_2})$  mean that  $k \in L^\infty(\mathbb{R}^3)$ .

*Remark 1.2.* The function  $k$ , which satisfies a Hölder condition of order  $\alpha$  with  $1 \leq \alpha < 3$  on  $H^1(\mathbb{R}^3)$  and achieves its maximum, is a special case of  $(H_{k_2})$ .

*Remark 1.3.* In Lemma 2.3, we show that  $\bar{\mu}$  is achieved.

By a solution  $(u, \phi)$  in  $H^1(\mathbb{R}^3) \times D^{1,2}(\mathbb{R}^3)$  of problem (1.1), we mean that for any  $v \in H^1(\mathbb{R}^3)$  it holds

$$\begin{cases} \int_{\mathbb{R}^3} (\nabla u \nabla v + uv + l(x)\phi uv) dx = \int_{\mathbb{R}^3} (k(x)|u|^{2^*-2}uv + \mu h(x)|u|^{q-2}uv) dx, \\ \int_{\mathbb{R}^3} \nabla \phi \nabla v dx = \int_{\mathbb{R}^3} l(x)u^2 v dx. \end{cases}$$

We say the solution is positive if  $u(x) > 0$  and  $\phi(x) > 0$  for all  $x \in \mathbb{R}^3$ .

We shall prove the following theorem.

**Theorem 1.4.** *Assume the hypotheses (H) hold and  $2 \leq q < 2^*$ . Then problem (1.1) has at least one positive solution  $(u, \phi_u)$  in  $H^1(\mathbb{R}^3) \times D^{1,2}(\mathbb{R}^3)$ .*

To prove the result above, we use a combination of techniques, e.g. techniques motivated by Willem [31], to overcome the lack of compactness of the Sobolev embedding, and methods used by Chen-Li-Li [11] and Zhao-Zhao [32], to estimate carefully the energy level.

**Notations.** Throughout this paper,  $L^p \equiv L^p(\mathbb{R}^3)$  ( $1 \leq p < +\infty$ ) is the usual Lebesgue space with the norm  $\|u\|_p^p = \int_{\mathbb{R}^3} |u|^p dx$ ;  $L^\infty \equiv L^\infty(\mathbb{R}^3)$  is the space of all essentially bounded functions with the norm  $\|u\|_\infty = \text{ess sup } |u|$ ;  $H^1 \equiv H^1(\mathbb{R}^3)$  denotes the usual Sobolev space with the norm  $\|u\|^2 = \int_{\mathbb{R}^3} (|\nabla u|^2 + |u|^2) dx$ ;  $H^{-1}$  is the dual space of  $H^1$  and  $\langle \cdot, \cdot \rangle \equiv \langle \cdot, \cdot \rangle_{H^{-1} \times H^1}$  is dual bracket;  $D^1 \equiv D^{1,2}(\mathbb{R}^3)$  is the completion of  $C_0^\infty(\mathbb{R}^3)$  with respect to the norm  $\|u\|_D^2 = \int_{\mathbb{R}^3} |\nabla u|^2 dx$ ;  $B_\rho(x)$  and  $B_\rho$  denote a ball with radius  $\rho$  centred at  $x$  and 0, respectively in a related space. Let  $u^+ = \max\{u, 0\}$  and  $u^- = \max\{-u, 0\}$ . We denote strong (weak) convergence for a sequence  $(u_n)_{n \in \mathbb{N}}$  and  $u$  in a Banach space by  $u_n \rightarrow u$  ( $u_n \rightharpoonup u$ ), respectively.  $N$  is used to denote the dimension, so  $N = 3$  if there is no special explanation. The so-called critical Sobolev exponent is denoted by  $2^* = \frac{2N}{N-2}$ . The symbol  $C$  denotes different positive constants and the value of  $C$  is allowed to change from line to line and in the same formula.

## 2 Preliminaries

In this section, we are going to give some preliminary lemmas. Since our methods are variational, first of all, it is necessary to transform the problem (1.1) into a Schrödinger

equation with a nonlocal term. In fact, for any  $u \in H^1$ , denote  $L_u(v)$  the linear functional in  $D^1$  by

$$L_u(v) = \int_{\mathbb{R}^3} l(x)u^2v dx.$$

It follows from the hypothesis  $(H_l)$ , Hölder and Sobolev inequalities that

$$|L_u(v)| \leq \|l\|_\infty \|u\|_{12/5}^2 \|v\|_6 \leq C \|l\|_\infty \|u\|_{12/5}^2 \|v\|_D. \quad (2.1)$$

Hence, the Lax-Milgram theorem implies that there exists, for each  $u$  in  $H^1$ , a unique  $\phi_u \in D^1$  such that

$$\int_{\mathbb{R}^3} \nabla \phi_u \nabla v = \int_{\mathbb{R}^3} l(x)u^2v dx \quad \text{for any } v \in D^1,$$

i.e.,  $\phi_u$  is the weak solution of  $-\Delta \phi = l(x)u^2$ . It holds

$$\phi_u(x) = \frac{1}{4\pi} \int_{\mathbb{R}^3} \frac{l(y)u^2(y)}{|x-y|} dy.$$

In particular, we have

$$\|\phi_u\|_D^2 = \int_{\mathbb{R}^3} |\nabla \phi_u|^2 dx = \int_{\mathbb{R}^3} l(x)\phi_u u^2 dx. \quad (2.2)$$

Using (2.1) and (2.2), we obtain

$$\|\phi_u\|_6 \leq C \|\phi_u\|_D \leq C \|u\|_{12/5}^2 \leq C \|u\|^2 \quad (2.3)$$

and

$$\int_{\mathbb{R}^3} l(x)\phi_u(x)u^2(x) dx \leq C \|u\|^4.$$

Thus  $F : H^1 \rightarrow \mathbb{R}$  is well defined with

$$F(u) = \int_{\mathbb{R}^3} l(x)\phi_u(x)u^2(x) dx. \quad (2.4)$$

To give the smoothness of the functional  $F$  (about the smoothness, we can find the statement in previous works, but we didn't find complete details), first, it is necessary to introduce the following lemma.

**Lemma 2.1.** [25, p.31] *Let  $0 < \beta < N$  and  $f \in L^q(\mathbb{R}^N)$ ,  $g \in L^r(\mathbb{R}^N)$  with  $\frac{1}{q} + \frac{1}{r} + \frac{\beta}{N} = 2$  and  $1 < q, r < \infty$ . Then*

$$\int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{|f(x)||g(y)|}{|x-y|^\beta} dx dy \leq C(q, r, \beta, N) \|f\|_q \|g\|_r, \quad x, y \in \mathbb{R}^N,$$

where  $C(q, r, \beta, N)$  is a positive constant depending on  $q, r, \beta$  and  $N$ .

**Lemma 2.2.** *If the hypothesis  $(H_l)$  holds, then  $F \in C^1(H^1, \mathbb{R})$ .*

*Proof.* From Lemma 2.1 and hypothesis  $(H_l)$  we obtain

$$\begin{aligned} & \int_{\mathbb{R}^3 \times \mathbb{R}^3} \frac{|l(x)u^2(x)||l(y)u(y)v(y)|}{|x-y|} dx dy \\ & \leq C \|u\|_{12/5}^2 \|uv\|_{6/5} \leq C \|u\|_{12/5}^2 \|u\|_{12/5} \|v\|_{12/5} \end{aligned}$$

for any  $u, v \in H^1$ . Then we may use the Lebesgue Theorem and Fubini Theorem and get

$$\begin{aligned} & \lim_{t \rightarrow 0} \frac{F(u+tv) - F(u)}{t} \\ & = \lim_{t \rightarrow 0} \int_{\mathbb{R}^3} \frac{l(x)}{t} \left( (u+tv)^2 \left( \phi_u + 2t \int_{\mathbb{R}^3} \frac{l(y)u(y)v(y)}{|x-y|} dy + t^2 \phi_v \right) - \phi_u u^2 \right) dx \\ & = 2 \int_{\mathbb{R}^3} l(x) \left( u^2(x) \int_{\mathbb{R}^3} \frac{l(y)u(y)v(y)}{|x-y|} dy + u(x)v(x) \int_{\mathbb{R}^3} \frac{l(y)u^2(y)}{|x-y|} dy \right) dx \\ & = 4 \int_{\mathbb{R}^3} l(x) \phi_u uv dx. \end{aligned}$$

Hence the Gateaux derivative of  $F$  on  $H^1$  exists and  $\langle \frac{1}{4} F'(u), v \rangle = \int_{\mathbb{R}^3} l(x) \phi_u uv dx$ . Let  $u_n \rightarrow u$  in  $H^1$  and  $v \in H^1$ , then by  $(H_l)$  we obtain

$$\begin{aligned} \|F'(u_n) - F'(u)\|_{H^{-1}} & = \sup_{\|v\|=1} |\langle F'(u_n) - F'(u), v \rangle| \\ & = 4 \sup_{\|v\|=1} \left| \int_{\mathbb{R}^3} l(x) (\phi_{u_n} u_n - \phi_{u_n} u + \phi_{u_n} u - \phi_u u) v dx \right| \\ & \leq 4 \|l\|_{\infty} \sup_{\|v\|=1} \left( \|\phi_{u_n}\|_6 \|u_n - u\|_{12/5} \|v\|_{12/5} + \int_{\mathbb{R}^3} |\phi_{u_n} - \phi_u| |uv| dx \right). \end{aligned} \quad (2.5)$$

It follows from Lemma 2.1 that

$$\begin{aligned} & \int_{\mathbb{R}^3} |\phi_{u_n} - \phi_u| |uv| dx \\ & = \int_{\mathbb{R}^3 \times \mathbb{R}^3} \frac{|u(x)v(x)||u_n^2(y) - u^2(y)|}{|x-y|} dx dy \\ & \leq C \|u_n^2 - u^2\|_{6/5} \|uv\|_{6/5} \leq C \|u_n^2 - u^2\|_{6/5} \|u\|_{12/5} \|v\|_{12/5}. \end{aligned}$$

From (2.3), (2.5), (2.6) and the fact that  $u_n \rightarrow u$  in  $H^1$ , we obtain

$$\|F'(u_n) - F'(u)\|_{H^{-1}} \rightarrow 0.$$

Thus  $F$  has a continuous Gateaux derivative on  $H^1$ . Therefore  $F \in C^1(H^1, \mathbb{R})$ .  $\square$

Let's introduce the Euler functional of the problem (1.1) as  $I : H^1 \rightarrow \mathbb{R}$  defined by

$$I(u) = \frac{1}{2} \|u\|^2 + \frac{1}{4} F(u) - \int_{\mathbb{R}^3} \left( \frac{1}{2^*} k(x) |u^+|^{2^*} + \frac{\mu}{q} h(x) |u^+|^q \right) dx. \quad (2.6)$$

By Lemma 2.2 we know that the functional  $I$  is of class  $C^1$  and its critical points are weak solutions of (1.1).

To prove Theorem 1.4, we still need some other preliminary lemmas.

**Lemma 2.3.** *Assume that the hypothesis  $(H_l)$  holds. Then  $F$  is a weakly continuous functional.*

*Proof.* Suppose  $u_n \rightharpoonup u$  in  $H^1$ . Since  $u_n \rightarrow u$  in  $L^2_{loc}$ , going if necessary to a subsequence, we can assume that

$$u_n \rightarrow u \text{ a.e. in } \mathbb{R}^3 \quad \text{and} \quad \phi_{u_n} \rightarrow \phi_u \text{ a.e. in } \mathbb{R}^3.$$

In fact, the last statement is true since, by  $(H_l)$  and Hölder inequality, we have

$$\begin{aligned} |\phi_{u_n}(x) - \phi_u(x)| &\leq \frac{1}{4\pi} \int_{\mathbb{R}^3} |l(y)| |u_n^2(y) - u^2(y)| \frac{1}{|x-y|} dy \\ &\leq C \|u_n^2 - u^2\|_{L^2(B_R(x))} \left( \int_{|x-y| \leq R} \frac{1}{|x-y|^2} dy \right)^{1/2} \\ &\quad + C \|u_n^2 - u^2\|_{L^{4/3}(B_R^c(x))} \left( \int_{|x-y| > R} \frac{1}{|x-y|^4} dy \right)^{1/4} \\ &\leq C \|u_n^2 - u^2\|_{L^2(B_R(x))} + CR^{-1/4} \|u_n^2 - u^2\|_{L^{4/3}(B_R^c(x))} \\ &\rightarrow 0, \end{aligned} \tag{2.7}$$

as  $n \rightarrow \infty$  and  $R \rightarrow \infty$ . Then  $\phi_{u_n} u_n^2 \rightarrow \phi_u u^2$  a.e. on  $\mathbb{R}^3$ . Moreover, the sequence  $(\phi_{u_n} u_n^2)_{n \in \mathbb{N}}$  is bounded in  $L^2$ , since

$$\int_{\mathbb{R}^3} (\phi_{u_n} u_n^2)^2 dx \leq \left( \int_{\mathbb{R}^3} \phi_{u_n}^6 dx \right)^{1/3} \left( \int_{\mathbb{R}^3} u_n^6 dx \right)^{2/3} = \|\phi_{u_n}\|_6^2 \|u_n\|_6^4 \leq C \|u_n\|_6^6.$$

Hence  $\phi_{u_n} u_n^2 \rightharpoonup \phi_u u^2$  in  $L^2$ . By  $(H_l)$  we have

$$F(u_n) = \int_{\mathbb{R}^3} l(x) \phi_{u_n} u_n^2 dx \rightarrow \int_{\mathbb{R}^3} l(x) \phi_u u^2 dx = F(u).$$

We have proved that  $F$  is weakly continuous. □

**Lemma 2.4.** *Assume the hypothesis  $(H_l)$  holds. Let  $u_n \rightharpoonup u$  in  $H^1$ , then*

$$F(u_n - u) = F(u_n) - F(u) + o(1).$$

*Proof.* Since  $(H_l)$  holds, from the proof of [32, Lemma 2.1], the result follows. □

From a similar proof as in [31, Lemma 2.13], we obtain the next result.

**Lemma 2.5.** *If the hypothesis  $(H_{h_1})$  holds and  $2 \leq q < 6$ , then the functional*

$$\psi_h : H^1 \rightarrow \mathbb{R} : u \mapsto \int_{\mathbb{R}^3} h(x) |u|^q dx$$

*is weakly continuous.*

**Lemma 2.6.** *Suppose the hypothesis  $(H_{h_1})$  holds and  $2 \leq q < 4$ . Then the following infimum*

$$\bar{\mu} := \mu_h = \inf_{u \in H^1 \setminus \{0\}} \left\{ \int_{\mathbb{R}^3} (|\nabla u|^2 + u^2) dx : \int_{\mathbb{R}^3} h(x) |u|^q dx = 1 \right\} \tag{2.8}$$

*is achieved.*

*Proof.* Let  $(u_n)_{n \in \mathbb{N}} \subset H^1$  be a minimizing sequence such that

$$\int_{\mathbb{R}^3} h(x)|u_n|^q dx = 1 \quad \text{and} \quad \int_{\mathbb{R}^3} (|\nabla u_n|^2 + u_n^2) dx \rightarrow \mu_h, \quad \text{as } n \rightarrow \infty.$$

So  $(u_n)_{n \in \mathbb{N}}$  is bounded in  $H^1$ . Then there exists a subsequence satisfying  $u_n \rightharpoonup u$  in  $H^1$ . Since  $h \in L^{6/(6-q)}$ , by Lemma 2.5, we have

$$\int_{\mathbb{R}^3} h(x)|u_n|^q dx \rightarrow \int_{\mathbb{R}^3} h(x)|u|^q dx. \quad \text{Hence} \quad \int_{\mathbb{R}^3} h(x)|u|^q dx = 1.$$

Then, by the weakly lower semi-continuous property of the norm, we get

$$\mu_h = \liminf_{n \rightarrow \infty} \int_{\mathbb{R}^3} (|\nabla u_n|^2 + u_n^2) dx \geq \int_{\mathbb{R}^3} (|\nabla u|^2 + u^2) dx \geq \mu_h.$$

Thus the infimum  $\mu_h$  is achieved.  $\square$

**Lemma 2.7.** *Suppose the hypotheses  $(H_l)$ ,  $(H_{k_1})$ ,  $(H_{h_1})$  and  $(H_{h_\mu})$  hold. Then  $I(0) = 0$  and  $(I_1)$  there are constants  $\rho, \alpha > 0$  such that  $I|_{\partial B_\rho} \geq \alpha$ ; and  $(I_2)$  there is  $\bar{u} \in H^1 \setminus \bar{B}_\rho$  such that  $I(\bar{u}) < 0$ .*

*Proof.* It is clear from the definition of  $I$  that  $I(0) = 0$ . To prove  $(I_1)$  and  $(I_2)$ , we consider  $2 \leq q < 4$  and  $4 \leq q < 6$  respectively. First, for  $2 \leq q < 4$ , we have  $0 < \mu < \bar{\mu}$  by  $(H_{h_\mu})$ . It follows from  $(H_{k_1})$ , Lemma 2.6 and Sobolev inequality that

$$\begin{aligned} I(u) &= \frac{1}{2}\|u\|^2 + \frac{1}{4}F(u) - \frac{1}{2^*} \int_{\mathbb{R}^3} k(x)|u^+|^{2^*} dx - \frac{\mu}{q} \int_{\mathbb{R}^3} h(x)|u^+|^q dx \\ &\geq \frac{1}{2}\|u\|^2 - C\|u\|^{2^*} - \frac{\mu}{q\bar{\mu}}\|u\|^2 = \|u\|^2 \left( \frac{1}{2} - \frac{\mu}{q\bar{\mu}} - C\|u\|^{2^*-2} \right). \end{aligned}$$

Set  $\rho = \|u\|$ , small enough such that  $C\rho^{2^*-2} \leq \frac{1}{2}(\frac{1}{2} - \frac{\mu}{q\bar{\mu}})$ . Hence we have

$$I(u) \geq \frac{1}{2} \left( \frac{1}{2} - \frac{\mu}{q\bar{\mu}} \right) \rho^2. \quad (2.9)$$

Take  $\alpha = \frac{1}{2}(\frac{1}{2} - \frac{\mu}{q\bar{\mu}})\rho^2$ . Then we get the result  $(I_1)$ . By (2.3) and the fact that  $\mu h(x) \geq 0$ , for fixed  $u_0$  with  $\|u_0\| = 1$  and  $\text{supp}(u_0) \subset \text{supp}(k)$ , we have

$$I(tu_0) \leq t^{2^*} \left( \frac{1}{2t^4}\|u_0\|^2 + \frac{C}{4t^2}\|u_0\|^4 - \frac{C}{2^*} \int_{\mathbb{R}^3} k(x)|u_0^+|^{2^*} dx \right).$$

Let  $t$  be large enough such that  $t > \rho$  and

$$\frac{1}{2t^4}\|u_0\|^2 + \frac{C}{4t^2}\|u_0\|^4 - \frac{C}{2^*} \int_{\mathbb{R}^3} k(x)|u_0^+|^{2^*} dx < 0.$$

Take  $\bar{u} = tu_0$ . Then  $(I_2)$  follows.

Next, we consider  $4 \leq q < 6$ , so  $\mu > 0$  by  $(H_{h_\mu})$ . Since  $(H_{k_1})$  and  $(H_{h_1})$  hold, the Hölder inequality and Sobolev inequality implies that

$$\begin{aligned} I(u) &= \frac{1}{2}\|u\|^2 + \frac{1}{4}F(u) - \frac{1}{2^*} \int_{\mathbb{R}^3} k(x)|u^+|^{2^*} dx - \frac{\mu}{q} \int_{\mathbb{R}^3} h(x)|u^+|^q dx \\ &\geq \frac{1}{2}\|u\|^2 - C\|u\|^{2^*} - \frac{\mu}{q} \|h\|_{\frac{6}{6-q}} \|u\|_6^q \\ &\geq \|u\|^2 \left( \frac{1}{2} - C\|u\|^{2^*-2} - C\|u\|^{q-2} \right) \end{aligned}$$

for each  $\mu > 0$  fixed. Hence  $(I_1)$  follows from the similar estimate with (2.9). The proof of  $(I_2)$  is the same to the case  $2 \leq q < 4$ .  $\square$

### 3 The proof of Theorem 1.4

To prove Theorem 1.4, we will apply the Mountain Pass Theorem to find a solution of problem (1.1) and then prove that it is a positive solution. Let us first recall (one of the versions of) the Mountain Pass Theorem.

**Mountain Pass Theorem [1].** Let  $E$  be a real Banach space and  $I \in C^1(E, \mathbb{R})$ . Suppose  $I(0) = 0$  and

$(I_1)$  there are constants  $\rho, \alpha > 0$  such that  $I|_{\partial B_\rho} \geq \alpha$ ; and

$(I_2)$  there is  $\bar{u} \in E \setminus \bar{B}_\rho$  such that  $I(\bar{u}) < 0$ . If  $I$  satisfies the  $(PS)_c$ -condition, where  $c$  is defined as

$$c = \inf_{g \in \Gamma} \max_{u \in g[0,1]} I(u) \quad \text{with } \Gamma = \{g \in C([0,1], E) : g(0) = 0, g(1) = \bar{u}\}. \quad (3.1)$$

Then  $I$  possesses a critical value  $c \geq \alpha$ .

Since Lemma 2.7 shows that the functional  $I$  has the Mountain Pass geometry, to apply this theorem to the functional  $I$  with  $E \equiv H^1$ , it is enough to prove that the Palais-Smale condition holds at the level  $c$  (the  $(PS)_c$ -condition for short), which means that every sequence  $(u_n)_{n \in \mathbb{N}} \subset H^1$  such that  $I(u_n) \rightarrow c$  and  $I'(u_n) \rightarrow 0$  in  $H^{-1}$  implies that  $(u_n)_{n \in \mathbb{N}}$  possesses a convergent subsequence in  $H^1$ .

**Lemma 3.1.** Assume  $(H_l)$ ,  $(H_{k_1})$ ,  $(H_{h_1})$  and  $(H_{h_\mu})$  hold. Then the functional  $I$  satisfies the  $(PS)_c$ -condition for  $c \in \left(0, \frac{1}{N} \mathcal{S}^{\frac{N}{2}} \|k\|_\infty^{-\frac{N-2}{2}}\right)$ , where  $\mathcal{S}$  denotes the best Sobolev constant defined by

$$\mathcal{S} = \inf_{u \in D^1 \setminus \{0\}} \frac{\int_{\mathbb{R}^3} |\nabla u|^2 dx}{\left(\int_{\mathbb{R}^3} |u|^{2^*} dx\right)^{2/2^*}}. \quad (3.2)$$

*Proof.* Let  $(u_n)_{n \in \mathbb{N}}$  be a  $(PS)_c$ -sequence of  $I$  at the level  $c \in \left(0, \frac{1}{N} \mathcal{S}^{\frac{N}{2}} \|k\|_\infty^{-\frac{N-2}{2}}\right)$ , i.e.,

$$I(u_n) \rightarrow c \quad \text{and} \quad I'(u_n) \rightarrow 0 \quad \text{in } H^{-1}. \quad (3.3)$$



**Step 1.** We consider  $2 \leq q < 4$ , so we get  $0 < \mu < \bar{\mu}$  by  $(H_{h_\mu})$ . Then by the Sobolev inequality, Lemma 2.6 and  $k(x) \geq 0$  for any  $x \in \mathbb{R}^3$ , for large  $n$  we have

$$\begin{aligned}
c + 1 + \|u_n\| &\geq I(u_n) - \frac{1}{4} \langle I'(u_n), u_n \rangle \\
&= \frac{1}{4} \|u_n\|^2 + \left( \frac{1}{4} - \frac{1}{2^*} \right) \int_{\mathbb{R}^3} k(x) |u_n^+|^{2^*} dx + \left( \frac{\mu}{4} - \frac{\mu}{q} \right) \int_{\mathbb{R}^3} h(x) |u_n^+|^q dx \\
&\geq \frac{1}{4} \|u_n\|^2 + \left( \frac{1}{4} - \frac{1}{2^*} \right) \int_{\mathbb{R}^3} k(x) |u_n^+|^{2^*} dx + \left( \frac{1}{4} - \frac{1}{q} \right) \frac{\mu}{\bar{\mu}} \|u_n\|^2 \\
&\geq \left( \frac{1}{4} + \left( \frac{1}{4} - \frac{1}{q} \right) \frac{\mu}{\bar{\mu}} \right) \|u_n\|^2,
\end{aligned} \tag{3.4}$$

which implies  $(u_n)_{n \in \mathbb{N}}$  is bounded in  $H^1$ , since  $0 < \mu < \bar{\mu}$  and  $2 \leq q < 4$ . Passing if necessary to a subsequence, we can assume that

$$\begin{aligned}
u_n &\rightharpoonup u \quad \text{in } H^1, \quad u_n \rightarrow u \quad \text{a.e. in } \mathbb{R}^3, \\
\nabla u_n &\rightharpoonup \nabla u \quad \text{in } L^2, \quad \text{and } u_n \rightarrow u \quad \text{in } L^2.
\end{aligned}$$

Let us define  $w_n = k(x) |u_n^+|^{N+2/N-2}$  and  $w = k(x) |u^+|^{N+2/N-2}$ . Since  $(u_n)_{n \in \mathbb{N}}$  is bounded in  $L^{2^*}$  and  $k \in L^\infty$ , then  $w_n$  is bounded in  $L^{2N/N+2}$  and so  $w_n \rightharpoonup w$  in  $L^{2N/N+2}$ . Note that for any  $v \in H^1$ , we have  $v \in L^{2N/N-2}$ ,  $\nabla v \in L^2$  and  $v \in L^2$ . Hence

$$\int_{\mathbb{R}^3} w_n v dx \rightarrow \int_{\mathbb{R}^3} w v dx, \quad \text{i.e.,} \quad \int_{\mathbb{R}^3} k(x) |u_n^+|^{2^*-1} v dx \rightarrow \int_{\mathbb{R}^3} k(x) |u^+|^{2^*-1} v dx, \tag{3.5}$$

and

$$\int_{\mathbb{R}^3} (\nabla u_n \nabla v + u_n v) dx \rightarrow \int_{\mathbb{R}^3} (\nabla u \nabla v + uv) dx. \tag{3.6}$$

From the proof of Lemma 2.3 and Lemma 2.5 we also have

$$\int_{\mathbb{R}^3} h(x) |u_n^+|^{q-1} v dx \rightarrow \int_{\mathbb{R}^3} h(x) |u^+|^{q-1} v dx, \tag{3.7}$$

and

$$\int_{\mathbb{R}^3} l(x) \phi_{u_n} u_n v dx \rightarrow \int_{\mathbb{R}^3} l(x) \phi_u u v dx. \tag{3.8}$$

Combining (3.5)–(3.8), for  $u_n \rightharpoonup u$  in  $H^1$ , we obtain

$$\begin{aligned}
\langle I'(u_n), v \rangle &= \int_{\mathbb{R}^3} (\nabla u_n \nabla v + u_n v) dx + \int_{\mathbb{R}^3} l(x) \phi_{u_n} u_n v dx \\
&\quad - \int_{\mathbb{R}^3} k(x) |u_n^+|^{2^*-1} v dx - \mu \int_{\mathbb{R}^3} h(x) |u_n^+|^{q-1} v dx \\
&\rightarrow \int_{\mathbb{R}^3} (\nabla u \nabla v + uv) dx + \int_{\mathbb{R}^3} l(x) \phi_u u v dx - \int_{\mathbb{R}^3} k(x) |u^+|^{2^*-1} v dx \\
&\quad - \mu \int_{\mathbb{R}^3} h(x) |u^+|^{q-1} v dx = \langle I'(u), v \rangle.
\end{aligned} \tag{3.9}$$

On the other hand, by the fact  $I'(u_n) \rightarrow 0$  in  $H^{-1}$ , we get that  $\langle I'(u_n), v \rangle \rightarrow 0$  for any  $v \in H^1$ . So  $\langle I'(u), v \rangle = 0$  for any  $v \in H^1$ , i.e.

$$-\Delta u + u + l(x) \phi_u u = k(x) |u^+|^{2^*-1} + \mu h(x) |u^+|^{q-1}. \tag{3.10}$$

In particular,  $\langle I'(u), u \rangle = 0$  and then from Lemma 2.6 and  $k(x) \geq 0$  we obtain

$$\begin{aligned} I(u) &= \frac{1}{4} \langle I'(u), u \rangle + \frac{1}{4} \|u\|^2 + \left( \frac{1}{4} - \frac{1}{2^*} \right) \int_{\mathbb{R}^3} k(x) |u^+|^{2^*} dx + \left( \frac{\mu}{4} - \frac{\mu}{q} \right) \int_{\mathbb{R}^3} h(x) |u^+|^q dx \\ &\geq \left( \frac{1}{4} + \left( \frac{1}{4} - \frac{1}{q} \right) \frac{\mu}{\bar{\mu}} \right) \|u\|^2 \geq 0. \end{aligned} \quad (3.11)$$

Let  $v_n = u_n - u$  and so  $v_n \rightarrow 0$  in  $H^1$ . Hence, using the given hypotheses, the Brézis-Lieb Lemma [7] implies that

$$\begin{aligned} \|u_n\|^2 &= \|v_n\|^2 + \|u\|^2 + o(1), \\ \int_{\mathbb{R}^3} k(x) |u_n^+|^{2^*} dx &= \int_{\mathbb{R}^3} k(x) |v_n^+|^{2^*} dx + \int_{\mathbb{R}^3} k(x) |u^+|^{2^*} dx + o(1), \\ \int_{\mathbb{R}^3} h(x) |u_n^+|^q dx &= \int_{\mathbb{R}^3} h(x) |v_n^+|^q dx + \int_{\mathbb{R}^3} h(x) |u^+|^q dx + o(1), \end{aligned}$$

and hence by Lemma 2.4 we have

$$I(u_n) = I(u) + \frac{1}{2} \|v_n\|^2 + \frac{1}{4} F(v_n) - \frac{1}{2^*} \int_{\mathbb{R}^3} k(x) |v_n^+|^{2^*} dx - \frac{1}{2} \int_{\mathbb{R}^3} h(x) |v_n^+|^q dx + o(1),$$

and

$$\langle I'(u_n), u_n \rangle = \langle I'(u), u \rangle + \|v_n\|^2 + F(v_n) - \int_{\mathbb{R}^3} k(x) |v_n^+|^{2^*} dx - \mu \int_{\mathbb{R}^3} h(x) |v_n^+|^q dx + o(1).$$

Therefore it follows from Lemma 2.3, Lemma 2.5 and the hypotheses  $I(u_n) \rightarrow c$  and  $I'(u_n) \rightarrow 0$  in  $H^{-1}$  that

$$c = \lim_{n \rightarrow \infty} I(u_n) = I(u) + \lim_{n \rightarrow \infty} \frac{1}{2} \|v_n\|^2 - \lim_{n \rightarrow \infty} \frac{1}{2^*} \int_{\mathbb{R}^3} k(x) |v_n^+|^{2^*} dx, \quad (3.12)$$

and

$$\langle I'(u), u \rangle + \lim_{n \rightarrow \infty} \|v_n\|^2 - \lim_{n \rightarrow \infty} \int_{\mathbb{R}^3} k(x) |v_n^+|^{2^*} dx = 0. \quad (3.13)$$

Using (3.10) and (3.13) we obtain

$$\|v_n\|^2 - \int_{\mathbb{R}^3} k(x) |v_n^+|^{2^*} dx \rightarrow -\langle I'(u), u \rangle = 0.$$

Now we may assume that

$$\|v_n\|^2 \rightarrow b \quad \text{and} \quad \int_{\mathbb{R}^3} k(x) |v_n^+|^{2^*} dx \rightarrow b.$$

By Sobolev's inequality we have

$$\|v_n\|^2 \geq \int_{\mathbb{R}^3} |\nabla v_n|^2 dx \geq \mathcal{S} \left( \int_{\mathbb{R}^3} |v_n^+|^{2^*} dx \right)^{2/2^*},$$

which means that

$$\int_{\mathbb{R}^3} k(x) |v_n^+|^{2^*} dx \leq \|k\|_\infty \int_{\mathbb{R}^3} |v_n^+|^{2^*} dx \leq \|k\|_\infty (\mathcal{S}^{-1} \|v_n\|^2)^{2^*/2},$$

i.e.,  $b \leq \|k\|_\infty (S^{-1}b)^{2^*/2}$ . So we get that  $b = 0$  or  $b \geq S^{\frac{N}{2}} \|k\|_\infty^{-\frac{N-2}{2}}$ . Assume  $b \geq S^{\frac{N}{2}} \|k\|_\infty^{-\frac{N-2}{2}}$ . Then combining (3.11) and (3.12), we obtain

$$c \geq \frac{1}{2}b - \frac{1}{2^*}b = \frac{1}{N}b \geq \frac{1}{N}S^{\frac{N}{2}} \|k\|_\infty^{-\frac{N-2}{2}},$$

which contradicts the fact that  $c < \frac{1}{N}S^{\frac{N}{2}} \|k\|_\infty^{-\frac{N-2}{2}}$ . Hence  $b = 0$ .

**Step 2.** For  $4 \leq q < 6$  and  $\mu > 0$ , we obtain that

$$\begin{aligned} c + 1 + \|u_n\| &\geq I(u_n) - \frac{1}{4}\langle I'(u_n), u_n \rangle \\ &= \frac{1}{4}\|u_n\|^2 + \left(\frac{1}{4} - \frac{1}{2^*}\right) \int_{\mathbb{R}^3} k(x)|u_n^+|^{2^*} dx + \left(\frac{\mu}{4} - \frac{\mu}{q}\right) \int_{\mathbb{R}^3} h(x)|u_n^+|^q dx \geq \frac{1}{4}\|u_n\|^2, \end{aligned}$$

which implies that  $(u_n)_{n \in \mathbb{N}}$  is bounded in  $H^1$ . To finish this step, we just need to replace (3.4) in Step 1 by the above inequality. The rest of the proof is similar to Step 1, so we omit it here.  $\square$

**Lemma 3.2.** *Suppose the hypotheses (H) hold. Then  $c < \frac{1}{N}S^{\frac{N}{2}} \|k\|_\infty^{-\frac{N-2}{2}}$ .*

*Proof.* The idea here is to find a path in  $\Gamma$  such that the maximum of the functional  $I$  at this path is strictly less than  $\frac{1}{N}S^{\frac{N}{2}} \|k\|_\infty^{-(N-2)/2}$ . To construct this path, we need the extremal function  $u_{\varepsilon, x_0}$  for the embedding  $D^1 \hookrightarrow L^6$ , where

$$u_{\varepsilon, x_0} = C \frac{\varepsilon^{1/4}}{(\varepsilon + |x - x_0|^2)^{1/2}}.$$

Here  $C$  is a normalizing constant and  $x_0$  is given in  $(H_{k_2})$ . Let  $\varphi \in C_0^\infty$  be such that  $0 \leq \varphi \leq 1$ ,  $\varphi|_{B_{R_2}} \equiv 1$  and  $\text{supp } \varphi \subset B_{2R_2}$  for some  $R_2 > 0$ . Set  $v_\varepsilon = \varphi u_{\varepsilon, x_0}$  and then  $v_\varepsilon \in H^1$  with  $v_\varepsilon(x) \geq 0$  for each  $x \in \mathbb{R}^3$ . The following asymptotic estimates hold if  $\varepsilon$  is small enough (see Brézis-Nirenberg [8]):

$$\|\nabla v_\varepsilon\|_2^2 = k_1 + O(\varepsilon^{\frac{1}{2}}), \quad \|v_\varepsilon\|_{2^*}^2 = k_2 + O(\varepsilon), \quad (3.14)$$

$$\|v_\varepsilon\|_s^s = \begin{cases} O(\varepsilon^{\frac{s}{4}}) & s \in [2, 3), \\ O(\varepsilon^{\frac{s}{4}} |\ln \varepsilon|) & s = 3, \\ O(\varepsilon^{\frac{6-s}{4}}) & s \in (3, 6), \end{cases} \quad (3.15)$$

with  $k_1/k_2 = \mathcal{S}$ , and  $2 \leq s < 2^*$ . We know the path  $tv_\varepsilon \in \Gamma$ . For the rest, we will prove

$$\max_{t \geq 0} I(tv_\varepsilon) < \frac{1}{N}S^{\frac{N}{2}} \|k\|_\infty^{-(N-2)/2} \quad (3.16)$$

for small  $\varepsilon$ . Since  $I(tv_\varepsilon) \rightarrow -\infty$  as  $t \rightarrow \infty$ , there exists  $t_\varepsilon > 0$  such that  $I(t_\varepsilon v_\varepsilon) = \max_{t \geq 0} I(tv_\varepsilon)$ . Also by Lemma 2.7,  $\max_{t \geq 0} I(tv_\varepsilon) \geq \alpha > 0$ . Then we have  $I(t_\varepsilon v_\varepsilon) \geq \alpha > 0$ . Thus from the continuity of  $I$ , we may assume that there exists some positive  $t_0$  such that  $t_\varepsilon \geq t_0 > 0$ .

Moreover from  $I(tv_\varepsilon) \rightarrow -\infty$  as  $t \rightarrow \infty$  and  $I(t_\varepsilon v_\varepsilon) \geq \alpha > 0$ , we get that there exists  $T_0$  such that  $t_\varepsilon \leq T_0$ . Hence  $t_0 \leq t_\varepsilon \leq T_0$ . Let  $I(t_\varepsilon v_\varepsilon) = A(\varepsilon) + B(\varepsilon) + C(\varepsilon)$ , where

$$A(\varepsilon) = \frac{t_\varepsilon^2}{2} \int_{\mathbb{R}^3} |\nabla v_\varepsilon|^2 dx - \frac{t_\varepsilon^{2^*}}{2^*} \int_{\mathbb{R}^3} k(x_0) |v_\varepsilon|^{2^*} dx,$$

$$B(\varepsilon) = \frac{t_\varepsilon^{2^*}}{2^*} \int_{\mathbb{R}^3} k(x_0) |v_\varepsilon|^{2^*} dx - \frac{t_\varepsilon^{2^*}}{2^*} \int_{\mathbb{R}^3} k(x) |v_\varepsilon|^{2^*} dx,$$

and

$$C(\varepsilon) = \frac{t_\varepsilon^2}{2} \int_{\mathbb{R}^3} |v_\varepsilon|^2 dx + \frac{t_\varepsilon^4}{4} F(v_\varepsilon) - \frac{t_\varepsilon^{2\mu}}{2} \int_{\mathbb{R}^3} h(x) |v_\varepsilon|^q dx,$$

since  $v_\varepsilon^+ = v_\varepsilon$ . First, we claim that

$$A(\varepsilon) \leq \frac{1}{N} S^{\frac{N}{2}} \|k\|_\infty^{-\frac{N-2}{2}} + C\varepsilon^{1/2}. \quad (3.17)$$

Indeed, let  $g(t) = \frac{t^2}{2} \int_{\mathbb{R}^3} |\nabla v_\varepsilon|^2 dx - \frac{t^{2^*}}{2^*} \int_{\mathbb{R}^3} k(x_0) |v_\varepsilon|^{2^*} dx$ . It is clear that  $g(t)$  achieves its maximum value at some  $T_\varepsilon$ . So

$$0 = g'(T_\varepsilon) = T_\varepsilon \int_{\mathbb{R}^3} |\nabla v_\varepsilon|^2 dx - T_\varepsilon^{2^*-1} \int_{\mathbb{R}^3} k(x_0) |v_\varepsilon|^{2^*} dx.$$

That is,

$$T_\varepsilon = \left( \frac{\int_{\mathbb{R}^3} |\nabla v_\varepsilon|^2 dx}{\int_{\mathbb{R}^3} k(x_0) |v_\varepsilon|^{2^*} dx} \right)^{\frac{1}{2^*-2}}.$$

Therefore, from (3.14), we have

$$g(T_\varepsilon) = \sup_{t \geq 0} g(t) = \frac{1}{N} \frac{\left( \int_{\mathbb{R}^3} |\nabla v_\varepsilon|^2 dx \right)^{N/2}}{\left( \int_{\mathbb{R}^3} k(x_0) |v_\varepsilon|^{2^*} dx \right)^{N-2/2}} = \frac{1}{N} S^{\frac{N}{2}} \|k\|_\infty^{-\frac{N-2}{2}} + C\varepsilon^{1/2}.$$

Then (3.17) follows. Secondly, we claim that  $B(\varepsilon) \leq C\varepsilon^{1/2}$ . In fact, since  $t_0 \leq t_\varepsilon \leq T_0$  and  $k \in L^\infty$ , by the definition of  $v_\varepsilon$ ,  $(H_{k_2})$  and using a change of variables with  $1 \leq \alpha < 3$ , we have

$$\begin{aligned} B(\varepsilon) &= \frac{t_\varepsilon^{2^*}}{2^*} \int_{\mathbb{R}^3} (k(x_0) - k(x)) |v_\varepsilon|^{2^*} dx \\ &\leq C\delta_1 \int_{|x-x_0| < \rho_1} \frac{|x-x_0|^\alpha \varepsilon^{3/2}}{(\varepsilon + |x-x_0|^2)^3} dx + C \int_{|x-x_0| \geq \rho_1} \frac{\varepsilon^{3/2}}{(\varepsilon + |x-x_0|^2)^3} dx \\ &\leq C\delta_1 \varepsilon^{\frac{3}{2}} \int_0^{\rho_1} \frac{r^{2+\alpha}}{(\varepsilon + r^2)^3} dr + C\varepsilon^{\frac{3}{2}} \int_{\rho_1}^\infty r^{-4} dr \\ &= C\delta_1 \varepsilon^{\frac{\alpha}{2}} \int_0^{\rho_1 \varepsilon^{-\frac{1}{2}}} \frac{\rho^{2+\alpha}}{(1 + \rho^2)^3} d\rho + C\rho_1^{-3} \varepsilon^{3/2} \\ &\leq C\delta_1 \varepsilon^{\frac{\alpha}{2}} + C\varepsilon^{3/2} \leq C\varepsilon^{\frac{3}{2}}. \end{aligned}$$

So we have proved our claim. Therefore, to finish the proof, it is enough to show

$$\lim_{\varepsilon \rightarrow 0^+} \frac{C(\varepsilon)}{\varepsilon^{1/2}} = -\infty. \quad (3.18)$$

Actually, from the definition of  $v_\varepsilon$ ,  $(H_{h_2})$  and for any  $\varepsilon$  such that  $0 < \varepsilon \leq \rho_2^2$ , it follows that

$$\begin{aligned} \int_{\mathbb{R}^3} h(x)|v_\varepsilon|^q dx &\geq C\delta_2 \int_{|x-x_0|<\rho_2} \frac{|x-x_0|^{-\beta}\varepsilon^{q/4}}{(\varepsilon+|x-x_0|^2)^{q/2}} dx + \int_{|x-x_0|\geq\rho_2} h(x)|v_\varepsilon|^q dx \\ &\geq C\delta_2\varepsilon^{q/4} \int_0^{\rho_2} \frac{r^2}{r^\beta(\varepsilon+r^2)^{q/2}} dr \\ &= C\delta_2\varepsilon^{\frac{3}{2}-\frac{q}{4}-\frac{\beta}{2}} \int_0^{\rho_2\varepsilon^{-\frac{1}{2}}} \frac{\rho^2}{\rho^\beta(1+\rho^2)^{q/2}} d\rho \\ &\geq C\delta_2\varepsilon^{\frac{3}{2}-\frac{q}{4}-\frac{\beta}{2}} \int_0^1 \frac{\rho^2}{2^q\rho^\beta} d\rho = C\varepsilon^{\frac{3}{2}-\frac{q}{4}-\frac{\beta}{2}}. \end{aligned}$$

Therefore, by the fact that  $t_0 \leq t_\varepsilon \leq T_0$  and hypothesis  $(H_I)$ , we have

$$\begin{aligned} C(\varepsilon) &= \frac{t_\varepsilon^2}{2} \int_{\mathbb{R}^3} |v_\varepsilon|^2 dx + \frac{t_\varepsilon^4}{4} F(v_\varepsilon) - \frac{t_\varepsilon^2\mu}{2} \int_{\mathbb{R}^3} h(x)|v_\varepsilon|^q dx \\ &\leq C\|v_\varepsilon\|_2^2 + C\|v_\varepsilon\|_{12/5}^4 - \mu C\varepsilon^{\frac{3}{2}-\frac{q}{4}-\frac{\beta}{2}} \\ &\leq C\varepsilon^{\frac{1}{2}} + C\varepsilon - \mu C\varepsilon^{\frac{3}{2}-\frac{q}{4}-\frac{\beta}{2}}. \end{aligned}$$

It follows from  $2 - \frac{q}{2} < \beta < 3$  that for fixed  $\mu$  we have

$$\frac{C(\varepsilon)}{\varepsilon^{1/2}} \leq C + C\varepsilon^{\frac{1}{2}} - \mu C\varepsilon^{1-\frac{q}{4}-\frac{\beta}{2}} \rightarrow -\infty, \text{ as } \varepsilon \rightarrow 0.$$

So we prove the claim (3.18). Therefore (3.16) follows.  $\square$

**Proof of Theorem 1.4.** It follows from Lemma 3.1 and Lemma 3.2 that the functional  $I$  satisfies the  $(PS)_c$ -condition at the level  $c$  defined by (3.1). And by Lemma 2.7, the functional  $I$  has the Mountain Pass geometry. Hence the functional  $I$  has a critical value  $c > 0$ . That is, there exists a nontrivial  $u \in H^1$  such that  $I'(u) = 0$ , which means that  $(u, \phi_u)$  is the nontrivial solution of system (1.1).

Since  $0 = \langle I'(u), u^- \rangle = \|u^-\|^2 + \int_{\mathbb{R}^3} l(x)\phi_u|u^-|^2 dx \geq \|u^-\|^2$ , then  $u \geq 0$  in  $\mathbb{R}^3$ . By standard arguments as in DiBenedetto [19] and Tolksdorf [28], we have that  $u \in L^\infty$  and  $u \in C_{loc}^{1,\gamma}$  with  $0 < \gamma < 1$ . Furthermore, by Harnack's inequality (see Trudinger [29]),  $u(x) > 0$  for any  $x \in \mathbb{R}^3$ . Thus  $(u, \phi_u)$  is a positive solution of system (1.1).  $\square$

## Acknowledgments

The authors thank the referees for their careful reading of the manuscript and insightful comments. This work was supported by FEDER funds through COMPETE – Operational Programme Factors of Competitiveness (“Programa Operacional Factores de Competitividade”) and by Portuguese funds through the *Center for Research and Development in Mathematics and Applications* (University of Aveiro) and the Portuguese Foundation for Science and Technology (“FCT–Fundaco para a Ciencia e a Tecnologia”), within project PEst-C/MAT/UI4106/2011 with COMPETE number FCOMP-01-0124-FEDER-022690. L. Huang also acknowledges the partial support of the PhD fellowship SFRH /BD/ 51162/2010.

## References

- [1] A. Ambrosetti and P. H. Rabinowitz, Dual variational methods in critical point theory and applications, *J. Funct. Anal.* **14** (1973), pp 349-381.
- [2] A. Ambrosetti and D. Ruiz, Multiple bound states for the Schrödinger-Poisson problem, *Commun. Contemp. Math.* **10** (2008), pp 391-404.
- [3] A. Azzollini and A. Pomponio, Ground state solutions for the nonlinear Schrödinger-Maxwell equations, *J. Math. Anal. Appl.* **345** (2008), pp 90-108.
- [4] V. Benci and D. Fortunato, An eigenvalue problem for the Schrödinger-Maxwell equations, *Topol. Methods Nonlinear Anal.* **11** (1998), pp 283-293.
- [5] V. Benci and D. Fortunato, Solitary waves of the nonlinear Klein-Gordon equation coupled with Maxwell equations, *Rev. Math. Phys.* **14** (2002), pp 409-420.
- [6] R. Benguria, H. Brézis and E. H. Lieb, The Thomas-Fermi-Von Weizsäcker theory of atoms and molecules, *Comm. Math. Phys.* **79** (1981), pp 167-180.
- [7] H. Brézis and E. H. Lieb, A relation between pointwise convergence of functions and convergence of functionals, *Proc. Amer. Math. Soc.* **8** (1983), pp 486-490.
- [8] H. Brézis and L. Nirenberg, Positive solutions of nonlinear elliptic equations involving critical Sobolev exponents, *Comm. Pure Appl. Math.* **36** (1983), pp 437-477.
- [9] I. Catto and P. L. Lions, Binding of atoms and stability of molecules in Hartree and Thomas-Fermi type theories. Part 1: A necessary and sufficient condition for the stability of general molecular system, *Comm. Partial Differential Equations*, **17** (1992), pp 1051-1110.
- [10] G. Cerami and G. Vaira, Positive solutions for some non-autonomous Schrödinger-Poisson systems, *J. Differential Equations*, **248** (2010), pp 521-543.
- [11] J. Chen, S. Li and Y. Li, Multiple solutions for a semilinear equation involving singular potential and critical exponent, *Z. angew. Math. Phys.* **56** (2005), pp 453-474.
- [12] G. M. Coclite, A Multiplicity result for the nonlinear Schrödinger-Maxwell equations, *Commun. Appl. Anal.* **7** (2003), pp 417-423.
- [13] T. D'Aprile and D. Mugnai, Solitary waves for nonlinear Klein-Gordon-Maxwell and Schrödinger-Maxwell equations, *Proc. Roy. Soc. Edinburgh Sect. A*, **134** (2004), pp 1-14.
- [14] T. D'Aprile and D. Mugnai, Non-existence results for the coupled Klein-Gordon-Maxwell equations, *Adv. Nonlinear Stud.* **4** (2004), pp 307-322.
- [15] T. D'Aprile and J. Wei, On bound states concentrating on spheres for the Maxwell-Schrödinger equations, *SIAM J. Math. Anal.* **37** (2005), pp 321-342.

- 
- [16] T. D'Aprile and J. Wei, Standing waves in the Maxwell-Schrödinger equations and an optimal configuration problem, *Calc. Var. Partial Differential Equations*, **25** (2006), pp 105-137.
- [17] P. D'Avenia, Non-radially symmetric solutions of nonlinear Schrödinger equation coupled with Maxwell equations, *Adv. Nonlinear Stud.* **2** (2002), pp 177-192.
- [18] P. D'Avenia, A. Pomponio and G. Vaira, Infinitely many positive solutions for a Schrödinger-Poisson system, *Nonlinear Anal.* **74** (2011), pp 5705-5721.
- [19] E. DiBenedetto,  $C^{1+\alpha}$  local regularity of weak solutions of degenerate elliptic equations, *Nonlinear Anal.* **7** (1983), pp 827-850.
- [20] X. He and W. Zou, Existence and concentration of ground states for Schrödinger-Poisson equations with critical growth, *J. Math. Phys.* **53** (2012), 023702.
- [21] G. Li, S. Peng and C. Wang, Multi-bump solutions for the nonlinear Schrödinger-Poisson system, *J. Math. Phys.* **52** (2011), 053505.
- [22] E. H. Lieb, Thomas-Fermi and related theories and molecules, *Rev. Modern Phys.* **53** (1981), pp 603-641.
- [23] P. L. Lions, Solutions of Hartree-Fock equations for Coulomb systems, *Comm. Math. Phys.* **109** (1984), pp 33-97.
- [24] P. Markowich, C. Ringhofer and C. Schmeiser, *Semiconductor Equations*, Springer-Verlag, New York, 1990.
- [25] M. Reed and B. Simon, *Methods of Modern Mathematical Physics*, Vols. II, IV, Elsevier (Singapore) Pte Ltd., 2003.
- [26] D. Ruiz, The Schrödinger-Poisson equation under the effect of a nonlinear local term, *J. Funct. Anal.* **237** (2006), pp 655-674.
- [27] J. Sun, H. Chen and J. J. Nieto, On ground state solutions for some non-autonomous Schrödinger-Poisson systems, *J. Differential Equations*, **252** (2012), pp 3365-3380.
- [28] P. Tolksdorf, Regularity for a more general class of quasilinear elliptic equations, *J. Differential Equations*, **51** (1984), pp 126-150.
- [29] N. S. Trudinger, On Harnack type inequality and their applications to quasilinear elliptic equations, *Comm. Pure Appl. Math.* **20** (1967), pp 721-747.
- [30] G. Vaira, Ground states for Schrödinger-Poisson type systems, *Ricerche mat.* **60** (2011), pp 263-297.
- [31] M. Willem, *Minimax Theorems*, Birkhäuser, Boston, 1996.
- [32] L. Zhao and F. Zhao, Positive solutions for Schrödinger-Poisson equations with a critical exponent, *Nonlinear Anal.* **70** (2009), pp 2150-2164.