## Supporting Information - A multi-scale model for the templated synthesis of mesoporous silica: The essential role of silica oligomers

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## Part I

## Molecular models

In this paper, we extend the coarse-grained MARTINI force field $\underline{\underline{1}}$ to systems containing silicate oligomers of different degrees of condensation. The MARTINI model is based on a mapping scheme whereby approximately four heavy (i.e., non-hydrogen) atoms are mapped onto one coarse-grained bead. As such, a single CG water bead represents four water molecules at the all-atom (AA) level. Details of the particular mapping used are provided below, together with tables of interaction parameters (Table S 11 and S 21 ). The MARTINI model has four

Table S 1 - Lennard-Jones parameter, $\varepsilon$, for the coarse-grained beads used in this work.

| epsilon | $\mathrm{Q}_{S I}$ | $\mathrm{P}_{4}$ | $\mathrm{BP}_{4}$ | $\mathrm{C}_{1}$ | $\mathrm{Q}_{0}$ | $\mathrm{Q}_{a}$ | $\mathrm{Q}_{d a}$ | $\mathrm{SQ}_{d a}$ | $\mathrm{SC}_{1}$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $\mathrm{Q}_{S I}$ | $5.6(\mathrm{O})$ | $4.5(\mathrm{II})$ | $4.5(\mathrm{II})$ | $3.5(\mathrm{IV})$ | $4.5(\mathrm{II})$ | $4.5(\mathrm{II})$ | $4.5(\mathrm{II})$ | $4.5(\mathrm{II})$ | $3.5(\mathrm{IV})$ |
| $\mathrm{P}_{4}$ | $4.5(\mathrm{II})$ | $5.0(\mathrm{I})$ | $5.6(\mathrm{O})$ | $2.0(\mathrm{VIII})$ | $5.6(\mathrm{O})$ | $5.6(\mathrm{O})$ | $5.6(\mathrm{O})$ | $5.6(\mathrm{O})$ | $2.0(\mathrm{VIII})$ |
| $\mathrm{BP}_{4}$ | $4.5(\mathrm{II})$ | $5.6(\mathrm{O})$ | $5.0(\mathrm{I})$ | $2.0(\mathrm{VIII})$ | $5.6(\mathrm{O})$ | $5.6(\mathrm{O})$ | $5.6(\mathrm{O})$ | $5.6(\mathrm{O})$ | $2.0(\mathrm{VIII})$ |
| $\mathrm{C}_{1}$ | $3.5(\mathrm{IV})$ | $2.0(\mathrm{VIII})$ | $2.0(\mathrm{VIII})$ | $5.6(\mathrm{O})$ | $2.0(\mathrm{IX})$ | $2.0(\mathrm{IX})$ | $2.0(\mathrm{IX})$ | $2.0(\mathrm{IX})$ | $3.5(\mathrm{IV})$ |
| $\mathrm{Q}_{0}$ | $4.5(\mathrm{II})$ | $5.6(\mathrm{O})$ | $5.6(\mathrm{O})$ | $2.0(\mathrm{IX})$ | $3.5(\mathrm{IV})$ | $4.5(\mathrm{II})$ | $4.5(\mathrm{II})$ | $4.5(\mathrm{II})$ | $2.0(\mathrm{IX})$ |
| $\mathrm{Q}_{a}$ | $4.5(\mathrm{II})$ | $5.6(\mathrm{O})$ | $5.6(\mathrm{O})$ | $2.0(\mathrm{IX})$ | $4.5(\mathrm{II})$ | $5.0(\mathrm{I})$ | $2.0(\mathrm{IX})$ | $2.0(\mathrm{IX})$ | $3.5(\mathrm{IV})$ |
| $\mathrm{Q}_{d a}$ | $4.5(\mathrm{II})$ | $5.6(\mathrm{O})$ | $5.6(\mathrm{O})$ | $2.0(\mathrm{IX})$ | $4.5(\mathrm{II})$ | $2.0(\mathrm{IX)}$ | $5.6(\mathrm{O})$ | $5.6(\mathrm{O})$ | $2.0(\mathrm{IX})$ |
| $\mathrm{SQ}_{d a}$ | $4.5(\mathrm{II})$ | $5.6(\mathrm{O})$ | $5.6(\mathrm{O})$ | $2.0(\mathrm{IX})$ | $4.5(\mathrm{II})$ | $2.0(\mathrm{IX})$ | $5.6(\mathrm{O})$ | 4.2 | $2.0(\mathrm{IX})$ |
| $\mathrm{SC}_{1}$ | $3.5(\mathrm{IV})$ | $2.0(\mathrm{IX})$ | $2.0(\mathrm{IX})$ | $3.5(\mathrm{IV})$ | $2.0(\mathrm{IX})$ | $3.5(\mathrm{IV})$ | $2.0(\mathrm{IX})$ | $2.0(\mathrm{IX})$ | 2.62 |

The capital letters in parentheses denote the interaction levels according to the standard nomenclature of the MARTINI force-field. Units of $\varepsilon$ are $\mathrm{kJ} / \mathrm{mol}$.

Table S $2-\sigma$ values for the coarse-grained beads used in this work. Units of $\sigma$ are in nm.

| sigma | $\mathrm{Q}_{S I}$ | $\mathrm{P}_{4}$ | $\mathrm{BP}_{4}$ | $\mathrm{C}_{1}$ | $\mathrm{Q}_{0}$ | $\mathrm{Q}_{a}$ | $\mathrm{Q}_{d a}$ | $\mathrm{SQ}_{d a}$ | $\mathrm{SC}_{1}$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| $\mathrm{Q}_{S I}$ | 0.47 | 0.47 | 0.47 | 0.47 | 0.47 | 0.47 | 0.47 | 0.47 | 0.47 |
| $\mathrm{P}_{4}$ | 0.47 | 0.47 | 0.47 | 0.47 | 0.47 | 0.47 | 0.47 | 0.47 | 0.47 |
| $\mathrm{BP}_{4}$ | 0.47 | 0.47 | 0.47 | 0.47 | 0.47 | 0.47 | 0.47 | 0.47 | 0.47 |
| $\mathrm{C}_{1}$ | 0.47 | 0.47 | 0.47 | 0.47 | 0.62 | 0.62 | 0.62 | 0.62 | 0.47 |
| $\mathrm{Q}_{0}$ | 0.47 | 0.47 | 0.47 | 0.62 | 0.47 | 0.47 | 0.47 | 0.47 | 0.62 |
| $\mathrm{Q}_{a}$ | 0.47 | 0.47 | 0.47 | 0.62 | 0.47 | 0.47 | 0.62 | 0.62 | 0.47 |
| $\mathrm{Q}_{d a}$ | 0.47 | 0.47 | 0.47 | 0.62 | 0.47 | 0.62 | 0.47 | 0.47 | 0.62 |
| $\mathrm{SQ}_{d a}$ | 0.47 | 0.47 | 0.47 | 0.62 | 0.47 | 0.62 | 0.47 | 0.43 | 0.62 |
| $\mathrm{SC}_{1}$ | 0.47 | 0.47 | 0.47 | 0.47 | 0.62 | 0.47 | 0.62 | 0.62 | 0.43 |

main types of interaction sites, namely polar (P), nonpolar (N), apolar (C) and charged (Q), with 18 subtypes overall. Q and N beads are subdivided taking into account their hydrogen-bonding capabilities: donor (d), acceptor (a), both (da) and none (0), whereas the degree of polarity in P and C beads is expressed by a number (from 1, low polarity, to 5, high polarity). A shifted Lennard-Jones (LJ) 12-6 potential energy function is used to describe the nonbonded interactions. The default effective bead size parameter is $\sigma=0.47$ nm . This is assumed for each interaction pair except for the two special classes of rings and antifreeze particles (see below), as well as for interactions between charged Q beads and the most apolar types $\left(\mathrm{C}_{1}\right)$ and $\left(\mathrm{C}_{2}\right)$ beads, where the range of repulsion is extended by setting $\sigma=0.62 \mathrm{~nm}$. The interaction strength is discretized in several levels, as follows: (O) 5.6; (I) 5.0; (II) 4.5; (III) 4.0; (IV) 3.5; (V) 3.1; (VI) 2.7; (VII); (VIII) 2.0 and (IX) 2.0 (with $\sigma=$ 0.62 nm ), all of them in $\mathrm{kJ} / \mathrm{mol}$ (more details can be consulted in ref. $\frac{1}{\text { ) }}$ ). In addition to the LJ interaction, charged groups (type Q) interact via a shifted Coulombic potential energy function with a relative dielectric constant of 15 . It should also be noted that the time scales of simulations using MARTINI should be rescaled by a factor of approximately four in order to reproduce realistic dynamic processes, because of the intrinsic speed-up caused by a reduction in the molecular degrees of freedom. $\underline{\underline{1}}$ However, in this paper we have chosen not to apply this correction, so the reported times are simply the simulation times.

An important feature of MARTINI is the mapping of ring particles by prefixed "S" beads with the aim to preserve the geometry of small ring compounds. In this set, the effective interaction size $\sigma$ is reduced to 0.43 rather than 0.47 nm , and the strength of ring-ring interactions $\varepsilon$ is reduced to $75 \%$ of the original value. This allows ring particles to pack more closely together without freezing in order to reproduce the experimental liquid densities of small ring compounds. $\underline{\underline{1}}$ Another challenging aspect of CG models is to accurately describe the behavior of water using a single LJ bead. In particular, a major concern is the unrealistic freezing of the CG model at temperatures above the melting temperature of real water. In
order to overcome this, Marrink et al. $\underline{\underline{1}}$ introduced antifreeze particles, $\left(\mathrm{BP}_{4}\right)$, which interact with all other particles in the system in exactly the same way as standard $\mathrm{P}_{4}$ water beads. However, the $\sigma$ value for interactions between $\mathrm{BP}_{4}$ and $\mathrm{P}_{4}$ is increased to 0.57 nm to disturb the lattice packing of the uniformly sized $\mathrm{P}_{4}$ particles. This essentially amounts to adding an entropic penalty to the freezing transition, and brings the model into better agreement with real water behavior.

Following the MARTINI philosophy, we mapped our $\mathrm{CTA}^{+}$surfactant molecule using five beads, four $\mathrm{C}_{1}$ for surfactant tails and $\mathrm{Q}_{0}$ for the head. TMA ${ }^{+}$and solvated $\mathrm{Br}^{-}$counterions were modeled by individual $\mathrm{Q}_{a}$ and $\mathrm{Q}_{0}$ beads, respectively. Benzene molecules were described as three connected $\mathrm{SC}_{1}$ beads, following the original MARTINI publication. $\underline{\underline{1}}$ The surfactant model was validated against experimental data and AA simulation in previous papers by ourselves $\underline{\underline{2}}^{\underline{2}}$ and others. $\underline{\underline{3}}$ Figure S 1 shows the mapping used in this paper for non-silicate molecules.


Figure S 1 - CG models for: (a) $\mathrm{CTA}^{+}$surfactant-four connected $\mathrm{C}_{1}$ beads for surfactant tails (hydrophobic) joined with a charged $\mathrm{Q}_{0}$ bead for the cationic surfactant head (hydrophilic) (b) one $\mathrm{Q}_{a}$ bead for bromide ions $\left(\mathrm{Br}^{-}\right)$with the first hydration shell (c) three interconnected $\mathrm{SC}_{1}$ beads for benzene (d) a single $\mathrm{Q}_{0}$ bead for tetramethylammonium (TMA ${ }^{+}$) cations (e) one bead for standard water molecules $\mathrm{P}_{4}$ or antifreeze $\mathrm{BP}_{4}$ particles, both with the $4: 1$ mapping.

To develop the CG force-field for silicates, we compared density profiles for both AA and CG models in preformed micelles containing silicates, to ensure that the correct physico-chemical behavior of each molecule was described at the CG level. The approach was described in detail in our previous paper, $\underline{\underline{2}}$ where we also presented the model for silica monomers as a single $\mathrm{Q}_{S I}$ bead. In this paper, we calibrated the parameters for dimers and higher oligomers using the same procedure. In the end, linear silica oligomers were represented by connected $\mathrm{Q}_{d a}$ beads (one for each Si-containing group), while cyclic silicates were described by $\mathrm{SQ}_{d a}$ beads (again, one per Si atom). Figures S 2 and S 3 show the mapping of our CG model for the silicate molecules used in this work.


Figure S 2 - CG models for linear silicates: (a) monomer with one $\mathrm{Q}_{S I}$ bead (b) two connected $\mathrm{Q}_{d a}$ beads for a dimer (c) three connected $\mathrm{Q}_{d a}$ beads for a linear trimer (d) a linear tetramer with four $\mathrm{Q}_{d a}$ connected beads.


Figure S 3 - CG models for cyclic silicates: (a) three interconnected $S Q_{d a}$ beads for a cyclic trimer (b) a cyclic tetramer with four $\mathrm{SQ}_{d a}$ interconnected beads (c) three interconnected $\mathrm{SQ}_{d a}$ beads joined with one $\mathrm{Q}_{d a}$ bead for branched cyclic tetramer (d) a double four ring (or cubic octamer) with eight $\mathrm{SQ}_{d a}$ interconnected beads.

The comparison of density profiles between AA and CG models for silica dimers was shown in the main paper, but Figure S 4 shows the (unsuccessful) comparison when the $\mathrm{Q}_{S I}$ potential parameters were transferred directly from monomers to dimers.


Figure S 4 - Density profile comparison, AA (dashed lines) and CG (solid lines), for the first attempt of silica dimers CG parameters using $Q_{S I}$ parameters. Color code is as follows: Surfactant tails in green, surfactant heads in purple, water in blue and silicates in red.

In Figure 5 5 whow the comparisons for linear silicates (trimer and tetramer), while in Figure $\sqrt{6} 6$ we show the comparisons for cyclic silicates (trimer, tetramer, branched cyclic trimer and cubic octamer). In all cases, we obtained good agreement between the AA and CG profiles.


Figure S 5 - Density profile comparison, AA (dashed lines) and CG (solid lines), for (a) linear trimers and (b) linear tetramers, using the $\mathrm{Q}_{d a}$ MARTINI parameters. Color code is the same as in Figure $\mathbb{S 4}$.


Figure S 6 - Density profile comparison, AA (dashed lines) and CG (solid lines), for (a) cyclic trimers, (b) cyclic tetramers, (c) branched cyclic trimers, and (d) double four ring octamers, using the $\mathrm{SQ}_{d a}$ MARTINI parameters. Color code is the same as in Figure S 4 .

## Part II

## Additional Results

This section includes additional figures that are relevant for the discussion presented in the main paper. Figure S 7 shows the detailed process of formation of a hexagonal mesophase starting from random distribution of silica dimers and surfactant. The system quickly formed small aggregates which fused into larger spherical micelles. No rod-like micelles were observed after this stage. Instead, micelles started to aggregate, changing their spherical shape into a more prolate shape. Several aggregates were formed in this way, which finally merged until the phase separation was complete. Finally, internal equilibration took place to yield an aggregate of small rod-like micelles similar to a HLC phase.


Figure S 7 - Diagram showing different stages of aggregation for a $6 \% \mathrm{w} / \mathrm{w}$ surfactant solution containing silicates in the form of $100 \%$ dimers, where all molecules were randomly placed at the beginning of the simulation. This corresponds to the results shown in Figure 2c of the main article. Color code is as follows: Surfactant tails in green, surfactant heads in purple, water in blue and silicates in red.

Figure S 8 shows a simulation snapshot for a surfactant/water system with a mixture of $50 \%$ monomers/dimers.


Figure S 8 - Snapshot showing a slice of the simulation box for the $50 \%$ monomer/dimer system. The slice is one-bead thick and the cutting plane was perpendicular to the main axis of the HLC mesostructure. Color code is as follows: Surfactant tails in green, surfactant heads in purple, water in blue, silica monomers in yellow and silica dimers in red.

We have cut out a slab with a thickness of one CG bead out of the HLC system to more clearly observe the core of the mesostructure. The figure shows the different role of silica monomers and dimers, with monomers well inside the rod surface whereas dimers are binding neighboring rods keeping the ordered structure. Several similar slices along the main axis of the HLC structure were taken, and in all of them the same pattern was found.

Firouzi et al. $\underline{4}$ showed that the addition of co-solvent molecules to an MCM-41 precursor solution induces a hexagonal-to-lamellar transition. Figure 89 shows the lamellar phase obtained after adding benzene molecules to a previously obtained hexagonal system (run 16) containing $\mathrm{CTA}^{+}$surfactants, water and silica dimers at 300 K . It is clear to see how the benzene molecules are mainly placed in the hydrophobic core of the lamellar phase (Figure S9b).

With the aim to estimate the contribution of adding benzenes in the hexagonal-to-lamellar

b)


Figure S 9 - Lamellar phase obtained after a hexagonal-to-lamellar transition promoted by adding benzene (run 16): (a) snapshot with all of the components except water molecules (b) snapshot showing only benzene molecules to highlight their location in the hydrophobic core of the lamellar phase. Color code is as follows: Hydrophobic surfactant tails in green, hydrophilic surfactant heads in purple, silica dimers in red and benzenes in brown.
transition (Figure S9), we have calculated the hydrophobic surfactant volume in both hexagonal (before adding benzene) and lamellar phases (after benzene addition). In case of the hexagonal phase, we estimated the hydrophobic core volume by measuring the diameter and length of each rod (Figure S10).

We took different hydrophobic core rod sections along the main axis to obtain an average


Figure S 10 - Rod diameters (a) and lengths (b) for the hexagonal system were visually obtained using the VMD $\underline{5}$ program. Only the external hydrophobic core beads in green color were plotted to facilitate the measurements. The system was split in different planes in order to enhance reading measurements.
diameter, and the hydrophobic core volume was obtained by multiplying each rod's crosssectional area by its length. This led to a total volume of $2343 \mathrm{~nm}^{3}$. In the case of the lamellar phase, we obtained the lamellae widths from the distance between the consecutive surfactant head density profile maximum. Then, to consider only the surfactant tail contribution we have subtracted the surfactant head bead size ( 0.47 nm ). Afterwards, the hydrophobic core volume was obtained by multiplying the above width by three (3 lamellae) and by the area of the lamellae (box size is 20.38 nm ). This yielded an approximate volume of $3564 \mathrm{~nm}^{3}$. Subtracting the hydrophobic core volume of the hexagonal phase from that of the lamellar phase, we can say that the lamellar system is $1221 \mathrm{~nm}^{3}$ larger than the hexagonal one. Values normalized by the total number of surfactant tails are provided in the main paper. On the other hand, the total volume of the added benzene molecules in the pure liquid phase was obtained by taking the density calculated by Marrink et al. $\underline{\underline{1}}$ for the benzene MARTINI model $\left(0.72 \mathrm{~g} / \mathrm{cm}^{3}\right)$. In our HLC system, we added 3000 benzene particles so the volume contribution was $540 \mathrm{~nm}^{3}$. This is significantly less than the effective volume increase during the transition.

Finally, Table S3 summarizes the number of components and the simulation box size which have been carried out in this work.

Table S 3 - Number of molecules and simulation box sizes at equilibrium for all the simulations carried out in this work.

|  | run 1 | run 2 | run 3 | run 4 | run 5 | run 6 | run 7 | run 8 | run 9 | run 10 | run 11 | run 12 | run 13 | run 14 | run 15 | run 16 | run 17 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{N}_{\text {Surf }}$ | 4000 | 4000 | 4000 | 4000 | 4000 | 2000 | 1800 | 1800 | 1000 | 999 | 1000 | 1000 | 1000 | 996 | 2000 | 4000 | 4000 |
| $\mathrm{N}_{B r}$ | 4000 | - | - | - | - | - | - | - | - | - | - | - | - | - | 2000 | - | - |
| $\mathrm{N}_{W}{ }^{*}$ | 320 K | 240K | 240K | 240K | 240K | 120K | 125K | 125K | 13K | 13K | 13K | 13K | 13K | 13K | 225K | 34.5K | 34.5K |
| $\mathrm{N}_{\text {SI1 }}$ | - | 4000 | - | - | 4000 | 1700 | 1200 | 600 | - | - | - | - | - | - | 256 | - | - |
| $\mathrm{N}_{\text {SI2 }}$ | - | - | 2000 | 2000 | - | 150 | 300 | 600 | 500 | - | - | - | - | - | - | 2000 | 2000 |
| $\mathrm{N}_{\text {SI3R }}$ | - | - | - | - | - | - | - | - | - | 333 | - | - | - | - | - | - | - |
| $\mathrm{N}_{\text {SI4R }}$ | - | - | - | - | - | - | - | - | - | - | 250 | - | - | - | - | - | - |
| $\mathrm{N}_{\text {D4R }}$ | - | - | - | - | - | - | - | - | - | - | - | 125 | 250 | 166 | 768 | - | - |
| $\mathrm{N}_{\text {TMA }}$ | - | - | - | - | - | - | - | - | - | - | - | - | - | - | 4864 | - | - |
| $\mathrm{N}_{B N Z}$ | - | - | - | - | - | - | - | - | - | - | - | - | - | - | - | 3000 | - |
| $\mathrm{L}_{x}$ | 35.61 | 32.58 | 32.52 | 32.50 | 32.58 | 25.85 | 26.08 | 26.07 | 8.44 | 8.44 | 8.45 | 8.36 | 8.41 | 8.20 | 31.74 | 20.38 | 19.78 |
| $\mathrm{L}_{y}$ | 35.61 | 32.58 | 32.52 | 32.50 | 32.58 | 25.85 | 26.08 | 26.07 | 8.44 | 8.44 | 8.45 | 8.36 | 8.41 | 8.20 | 31.74 | 20.38 | 19.78 |
| $\mathrm{L}_{z}$ | 35.61 | 32.58 | 32.52 | 32.50 | 32.58 | 25.85 | 26.08 | 26.07 | 32.72 | 32.72 | 32.75 | 33.45 | 33.63 | 34.87 | 31.74 | 20.38 | 19.78 |

number of silica dimers; $\mathrm{N}_{S I 3 R}$ number of silica ring trimers; $\mathrm{N}_{S I 4 R}$ number of silica ring tetramers; $\mathrm{N}_{D 4 R}$ number of silica cubic octamers; $\mathrm{N}_{T M A}$ number of tetramethylammonium cations; $\mathrm{N}_{B N Z}$ number of benzene molecules; $\mathrm{L}_{x}, \mathrm{~L}_{y}$ and $\mathrm{L}_{z}$ box size in nanometers for $\mathrm{x}, \mathrm{y}$ and z axis, respectively.

* K means $10^{3}$.


## References

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